

$$\Rightarrow d\mu \rightsquigarrow \frac{2^{r_1+r_2+1}}{\sqrt{r_1+r_2}} \cdot R \frac{dy_0}{y_0} dy_1 \dots dy_{r_1+r_2-1} \prod_{\sigma: \mathbb{R}} d\tau_\sigma \prod_{z: \mathbb{C}} dz_z =: d\omega$$

det  $(\text{wpt}(\Sigma_g^{(R)}))$

lem (i)  
(Boundary of  $\partial D_T$  in  $\gamma$ 's).

two parts: part of the bdy of  $\partial \Sigma_{\mathbb{O}_F \otimes \mathbb{O}_A}$  in pair  
 } hyperplane of  $\partial D_T$  defined by the distance  $d(\eta, z_1) = \frac{1}{T}$

becomes:  $\gamma_0 = T' = N(\text{corb}^{-1}) \cdot T$ .

(ii)  $d\mu|_{\gamma_0=T'} = (d\mu) \frac{\sqrt{r_1+r_2}}{T'} \cdot 2^{r_1+r_2+1} \cdot R dy_1 \dots dy_{r_1+r_2-1} \prod_{\sigma: \mathbb{R}} d\tau_\sigma \prod_{z: \mathbb{C}} dz_z$

(iii) of  $v$  is the unit normal to the bdy hyperplane of distance

$v = (\frac{1}{\sqrt{r_1+r_2} T'}, 0, \dots, 0)$  &  $\frac{\partial f}{\partial v} = (\frac{1}{\sqrt{r_1+r_2} T'}, 0, \dots, 0)$  grad  $f$

since  $\langle \frac{\partial}{\partial \gamma_0}, \frac{\partial}{\partial \gamma_0} \rangle = \frac{1}{\gamma_0^2} = \frac{1}{(r_1+r_2) T'^2} = \sqrt{r_1+r_2} \cdot T' \cdot \frac{\partial}{\partial \gamma_0}$

$$\Rightarrow \int_{\partial D_T} \hat{E}_{\text{vac}}(z, s) d\mu = \left( \frac{1}{r_1+r_2} - \frac{1}{s \text{corb}} \right) \sum_{i=1}^h \int_{X_i(T_i)} \frac{\partial \hat{E}_{\text{vac}}(z, s)}{\partial v} dS$$

they are in pair opposite direction.

$\gamma_0 = T'_i, i=1, \dots, h$

$$= \frac{r_1+r_2}{r_1+r_2} \cdot 2^{r_1+r_2+1} \cdot R \cdot (D_{PF})^{\frac{1}{2}} \cdot \sum_{i=1}^h N(\text{corb}_i^{-1}) \cdot \left( \frac{A_{0i}}{s_1} (T'_i)^{s_1} - \frac{B_{0i}}{s} (T'_i)^s \right)$$

Thus (Weyl) Up to a constant factor depending on  $F$ , we have

$$\int_{\partial D_T} \hat{E}_{\text{vac}}(z, s) d\mu = \frac{\hat{\Sigma}_P(z, s)}{s_1} \cdot T^{s_1} - \frac{\hat{\Sigma}_F(z, s)}{s} \cdot T^s$$